

Searching for H₂ emission from protoplanetary disks using near- and mid-infrared high-resolution spectroscopy

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Abstract. The mass and dynamics of protoplanetary disks are dominated by molecular hydrogen (H₂). However, observationally very little is known about the H₂. In this paper, we discuss two projects aimed to constrain the properties of H₂ in the disk's planet forming region (R < 50 AU). First, we present a sensitive survey for pure-rotational H₂ emission at 12.278 and 17.035 μ m in a sample of nearby Herbig Ae/Be and T Tauri stars using VISIR, ESO's VLT high-resolution mid-infrared spectrograph. Second, we report on a search for H₂ ro-vibrational emission at 2.1228, 2.2233 and 2.2477 μ m in the classical T Tauri star LkH α 264 and the debris disk 49 Cet employing CRIRES, ESO's VLT high-resolution near-infrared spectrograph.

VISIR project: none of the sources show H₂ mid-IR emission. The observed disks contain less than a few tenths of M_{Jupiter} of optically thin H₂ at 150 K, and less than a few M_{Earth} at $T > 300$ K. Our non-detections are consistent with the low flux levels expected from the small amount of H₂ gas in the surface layer of a Chiang and Goldreich (1997) Herbig Ae two-layer disk model. In our sources the H₂ and dust in the surface layer have not significantly departed from thermal coupling ($T_{\text{gas}}/T_{\text{dust}} < 2$) and the gas-to-dust ratio in the surface layer is very likely < 1000 .

CRILES project: The H₂ lines at 2.1218 μ m and 2.2233 μ m are detected in LkH α 264. An upper limit on the 2.2477 μ m H₂ line flux in LkH α 264 is derived. 49 Cet does not exhibit H₂ emission in any of observed lines. There are a few M_{Moon} of optically thin hot H₂ in the inner disk (~ 0.1 AU) of LkH α 264, and less than a tenth of a M_{Moon} of hot H₂ in the inner disk of 49 Cet. The shape of the 1–0 S(0) line indicates that LkH α disk is close to face-on ($i < 35^\circ$). The measured 1–0 S(0)/1–0 S(1) and 2–1 S(1)/1–0 S(1) line ratios in LkH α 264 indicate that the H₂ is thermally excited at $T < 1500$ K. The lack of H₂ emission in the NIR spectra of 49 Cet and the absence of H α emission suggest that the gas in the inner disk of 49 Cet has dissipated.

† This paper is based on materials of our A&A papers Carmona *et al.* 2008 & 2007.

Keywords. stars: emission-line – stars: pre-main sequence – planetary systems:protoplanetary disks

1. Introduction

Circumstellar disks surrounding low- and intermediate- mass stars in their pre-main sequence phase are the locations where planets presumably form. Such protoplanetary disks are composed of gas and dust. Their mass and dynamics are dominated by gas (99%), specifically by molecular hydrogen (H_2). Observationally, very little is known about the gas compared to the dust. However, if we want to answer major questions in planet formation such as: how massive are the disks?, how extended are the disks?, and how long protoplanetary disks last ? we require information about the gaseous component of the disk. In particular, we would like to characterize warm gas in the inner disk ($R < 50$ AU), the region where planets form (see reviews by Najita *et al.* 2007; Carr 2005).

H_2 is by far the most abundant molecular species in protoplanetary disks. Unfortunately, H_2 is one of the most challenging molecules to detect. H_2 is a homonuclear molecule that lacks a permanent dipole moment, therefore, H_2 transitions are electric quadrupole in nature, thus, very weak (i.e., small Einstein coefficients). In addition, in the case of protoplanetary disks, the H_2 lines are not sensitive to the gas in the optically thick regions where the dust and gas are at equal temperature. Practical observational challenges also have to be faced. The H_2 lines from the disk needs to be detected on the top of a strong infrared continuum. From the ground, the infrared windows (specially the mid-infrared) are strongly affected by sky and instrument background emission, and the H_2 transitions lie close to atmospheric absorption lines highly dependent on atmospheric conditions. The advent of high spectral resolution infrared spectrographs mounted on large aperture telescopes, allows for the first time the study of H_2 emission from the ground.

In this paper, we discuss two projects aimed to constrain the properties of H_2 gas in the planet forming region of circumstellar disks. In Sect. 1 we present the results of a large, sensitive survey for pure-rotational molecular hydrogen emission at 12.278 and 17.035 microns in a sample of nearby Herbig Ae/Be and T Tauri stars using the high-resolution mid-infrared spectrograph VISIR at ESO-VLT (Carmona *et al.* 2008). In Sect. 2. we report on the results of a search for H_2 ro-vibrational emission at 2.1228, 2.2233 and 2.2477 μm in the classical T Tauri star LkH α 264 and the debris disk 49 Ceti using CRIRES, the new ESO's VLT Adaptive Optics high resolution near-infrared spectrograph (Carmona *et al.* 2007).

2. Searching for mid-IR H_2 emission from protoplanetary disks with VISIR

2.1. Motivation.

To probe the gas in the giant planet forming region of the disk. H_2 mid-IR lines probe warm gas at $T \sim 150\text{--}1000$ K. This gas is located from a few AU up to 50 AU.

2.2. Previous work.

H_2 mid-IR emission from protoplanetary disks has been reported from ISO observations (Thi *et al.* 2001). However, subsequent ground-based efforts (Richter *et al.* 2002; Sheret *et al.* 2003; Sako *et al.* 2005) did not confirm the ISO detections. H_2 emission in the mid-IR has been searched towards debris disks using Spitzer (Hollenbach *et al.* 2005, Pascucci *et al.* 2006, Chen *et al.* 2006) with no detection reported. Most recently, Bitner *et al.* (2007) and Martin-Zaïdi *et al.* (2007) reported the detection of mid-IR H_2 emission in two Herbig Ae/Be stars (AB Aur and HD 97048) from the ground, and Lahuis *et al.* (2007) reported the detection of mid-IR H_2 emission in 6 T Tauri stars with *Spitzer*.

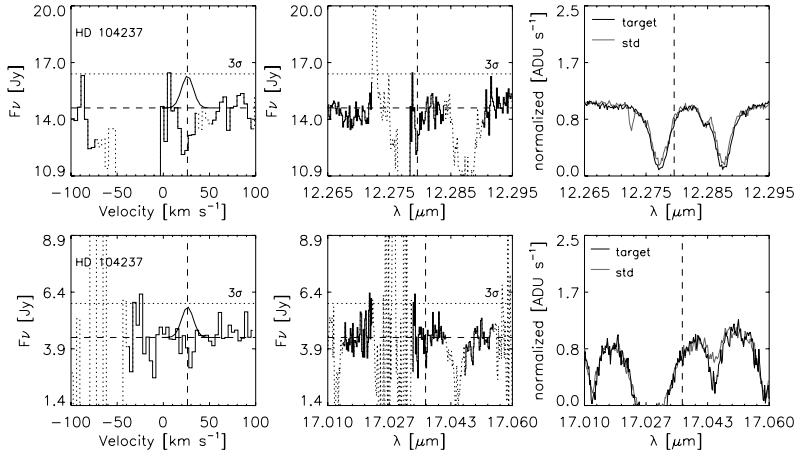


Figure 1. Example of the VISIR spectra obtained in the case of HD 104237. The upper panels show the spectra for the H₂ 0–0 S(2) line at 12.278 μm . The lower panels display the spectra for the H₂ 0–0 S(1) line at 17.035 μm . The left panels show a zoom to the -100 to 100 km s^{-1} interval of the atmospheric corrected spectra. A Gaussian of $FWHM = 15 \text{ km s}^{-1}$ and integrated line flux equal to the line-flux upper limits obtained is overplotted at the expected velocity shifted location (vertical dashed line). The central panels show the full corrected spectra. Dotted lines show spectral regions strongly affected by telluric or standard star absorption features. The right panels show the continuum normalized spectra of the standard star and the target before telluric correction. The spectra are not corrected for the radial velocity of the targets (Carmona *et al.* 2008).

2.3. Observations & Data Reduction.

We observed the Herbig Ae/Be stars UX Ori, HD 34282, HD 100453, HD 101412, HD 104237 and HD 142666, and the T Tauri star HD 319139 in the first semester of 2006 and 2007 with VISIR, a combined imager and spectrograph designed for observations in the N ($\approx 8 - 13 \mu\text{m}$) and Q bands ($\approx 16.5 - 24.5 \mu\text{m}$) (Lagage *et al.* 2004), mounted at the ESO-VLT Melipal telescope in Cerro Paranal, Chile. We selected a sample of well known nearby Herbig Ae/Be and T Tauri stars based on evidence of large disk reservoirs. The targets have either reported detections of cold CO gas at (sub)-mm wavelengths, or dust continuum emission at mm wavelengths. We chose stars with $12 \mu\text{m}$ continuum fluxes $> 0.5 \text{ Jy}$ (otherwise too faint for acquisition with VISIR) and $< 25 \text{ Jy}$ (hard to detect weak lines on top of a strong continuum). The H₂ $v = 0 - 0 S(1)$ line at $17.035 \mu\text{m}$ was observed in the *high-resolution long-slit mode* with a 0.4 arcsec slit, giving a spectral resolution $R \approx 21000$, or 14 km s^{-1} . The H₂ $v = 0 - 0 S(2)$ line at $12.278 \mu\text{m}$ was observed in the *high-resolution echelle mode* with a 0.4 arcsec slit, giving a spectral resolution $R \approx 20000$, or 15 km s^{-1} . The total integration time in each line was 1 h. The slit was oriented in the North-South direction. Sky background was subtracted by chopping the telescope by $\sim 8''$ in the direction of the slit. Asymmetrical thermal background of the telescope was subtracted by nodding the telescope by $\sim 8''$ in the direction of the slit.

For correcting the spectra for telluric absorption and obtaining the absolute flux calibration, spectroscopic standard stars at close airmasses to that of the science targets were observed immediately preceding and following the $12 \mu\text{m}$ exposure, and preceding or following the $17 \mu\text{m}$ exposure. After assuring that all the half-chop cycles in the VISIR data cubes had the same wavelength in the same row of pixels, the data-cubes were processed with the VISIR pipeline (Lundin 2006). The science spectrum was extracted by summing the number of counts inside the PSF in the dispersion direction in the 2D spectrum. To correct for telluric absorption and flux-calibrate the science spectrum, the one-dimensional extracted science spectrum was divided by the one-dimensional extracted

Table 1. H₂ mid-IR emission line flux upper limits and optically thin H₂ gas mass limits.

Star	H ₂ line	λ	continuum	line flux ^a	H ₂ mass limits in M _J		
		[μ m]	[Jy]	[$\times 10^{-14}$ ergs s ⁻¹ cm ⁻²]	$T = 150$ K	$T = 300$ K	$T = 1000$ K
UX Ori	0-0 S(2)	12.278	1.9 (1.1)	<1.4	27.9	1.9×10^{-1}	1.2×10^{-2}
	0-0 S(1)	17.035	2.0 (1.5)	<1.3	1.0	6.8×10^{-2}	2.0×10^{-2}
HD 34282	0-0 S(2)	12.278	0.3 (0.4)	<0.5	13.8	9.7×10^{-2}	6.0×10^{-3}
	0-0 S(1)	17.035
HD 100453	0-0 S(2)	12.278	9.0 (0.7)	<0.9	1.9	1.4×10^{-2}	0.9×10^{-3}
	0-0 S(1)	17.035	14.8 (1.7)	<1.5	0.1	0.9×10^{-2}	2.5×10^{-3}
HD 101412	0-0 S(2)	12.278	3.5 (1.0)	<1.2	5.3	3.7×10^{-2}	2.3×10^{-3}
	0-0 S(1)	17.035	1.6 (1.6)	<1.4	0.2	1.6×10^{-2}	4.8×10^{-3}
HD 104237	0-0 S(2)	12.278	14.6 (1.8)	<2.2	5.0	3.6×10^{-2}	2.2×10^{-3}
	0-0 S(1)	17.035	4.4 (1.5)	<1.3	0.1	0.8×10^{-2}	2.3×10^{-3}
HD 142666	0-0 S(2)	12.278
	0-0 S(1)	17.035	8.0 (1.2)	<1.1	0.2	1.0×10^{-2}	3.1×10^{-3}
HD 319139	0-0 S(2)	12.278
	0-0 S(1)	17.035	2.5 (1.5)	<1.3	0.2	1.2×10^{-2}	3.7×10^{-3}

Notes: ^a Upper limits calculated using a line *FWHM* of 15 km s⁻¹.

spectrum of the standard star (for further details on the data reduction procedure, telluric correction and flux calibration see Sect. 3 of Carmona *et al.* 2008).

2.4. Results.

None of the observed sources show evidence for H₂ emission at 12 or 17 μ m (see Fig. 1 for an example of the spectra obtained in the case of HD 104237, for all the spectra see Carmona *et al.* 2008). Assuming that the H₂ lines are unresolved, we calculated 3σ upper limits to the integrated H₂ line fluxes by multiplying the 3σ continuum flux noise with the instrument resolution line width (~ 15 km s⁻¹). Our results are summarized in Table 1. The typical sensitivity limit of our observations is a line flux of 10^{-14} ergs s⁻¹ cm⁻². Our flux limits are of the order of magnitude of the H₂ 0-0 S(1) and 0-0 S(2) lines fluxes (1.1 and 0.53×10^{-14} ergs s⁻¹ cm⁻², respectively) reported by Bitner *et al.* (2007) for AB Aur[†] and the H₂ line fluxes of 0.33 to 1.70×10^{-14} ergs s⁻¹ cm⁻² reported for the H₂ 0-0 S(2) line in the observations by Lahuis *et al.* (2007).

2.5. Discussion.

Under the assumption that the H₂ emission and the emission of the accompanying dust are optically thin, that the emitting H₂ is in local thermodynamical equilibrium (LTE), and that the source size is equal or smaller than VISIR's beam size, we derived upper limits to the H₂ mass as a function of the temperature employing (Thi *et al.* 2001)

$$M_{\text{gas}} = f \times 1.76 \times 10^{-20} \frac{4\pi d^2 F_{ul}}{E_{ul} A_{ul} x_u(T)} M_{\odot}, \quad (2.1)$$

where F_{ul} is the upper limit to the integrated line flux, d is the distance in pc to the star, E_{ul} is the energy of the transition, A_{ul} is the Einstein coefficient of the $J = u - l$ transition and x_u is the population of the level u at the excitation temperature T in LTE; f is the conversion factor required for deriving the total gas mass from the H₂-ortho or H₂-para mass determined. In Table 1, we present our results. The disks contain less than a few tenths of Jupiter mass of optically thin H₂ at 150 K, and less than a few Earth masses of optically thin H₂ at 300 K and higher temperatures.

Using a two-layer Chiang and Goldreich (1997, CG 97) disk model implementation (CGplus, Dullemond *et al.* 2001) with physical parameters aimed to fit the spectral

[†] Note that the observations of AB Aur were performed with TEXES, which provides a spectral resolution of 100000, increasing the line-to-continuum contrast compared with our observations.

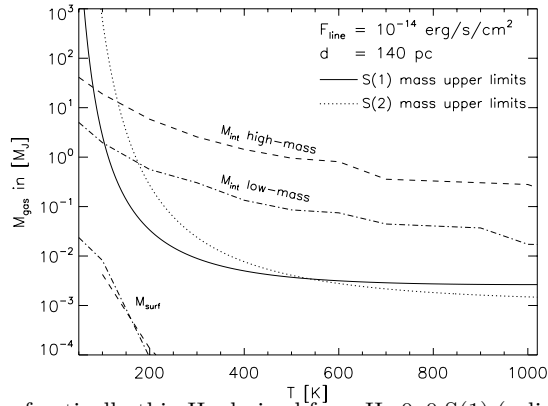


Figure 2. Mass limits of optically thin H₂ derived from H₂ 0–0 S(1) (solid line) and H₂ 0–0 S(2) (dotted line) as a function of the temperature for a line flux limit of 10^{-14} erg s^{−1} cm^{−2} for a source at a distance of 140 pc. Dashed and dot-dashed lines show the gas mass as function of the temperature for a Chiang and Goldreich (1997) optically thick two-layer model for a low-mass ($M_{\text{DISK}} = 0.02 M_{\odot}$) and a high-mass ($M_{\text{DISK}} = 0.11 M_{\odot}$) disk assuming a gas-to-dust ratio of 100. M_{int} is the mass of the interior layer. M_{surf} is the mass of the surface layer. H₂ emission arises only from the optically thin molecular gas in surface layer of the disk (Carmona *et al.* 2008).

energy distribution (SED) of prototypical Herbig Ae/Be stars (Chiang *et al.* 2001), we computed the expected amount of gas in the interior and surface layer as function of the temperature for disks of mass $0.02 M_{\odot}$ and $0.11 M_{\odot}$ (see Fig. 2). The mass limits derived from the H₂ 0–0 S(1) and 0–0 S(2) line observations are smaller than the amount of warm gas in the interior layer, but much larger than the amount of molecular gas in the surface layer. Fig. 2 shows that the amount of gas in the surface layer is very small ($<10^{-2} M_J \sim 3M_{\oplus}$) and almost independent of the total mass of the disk. *If the two-layer model is an adequate representation of the structure of the disk, the thermal flux levels of H₂ mid-infrared emission are below the detection limit of the observations, because the mass of H₂ in the surface layer is very small.*

In Carmona *et al.* (2008), based on a CG97 two-layer model, assuming LTE thermal emission, we calculated the H₂ mid-IR fluxes expected from optically thick Herbig Ae disks (see Sect. 5.2 of Carmona *et al.* 2008 for a detailed description of the calculation). We found that for a source distant 140 pc, the expected fluxes are of the order of 10^{-16} erg s^{−1} cm^{−2} for the 0–0 S(1) line and 10^{-17} erg s^{−1} cm^{−2} for the 0–0 S(2) line, and line-to continuum flux ratios of $<10^{-3}$. *These line flux levels are two orders of magnitude below the sensitivity limits of our observations (5×10^{-15} erg s^{−1} cm^{−2}).* If the two-layer approximation to the structure of the disk is correct, we are essentially “blind” to most of the warm H₂ in the disk because it is located in the optically thick interior layer of the disk.

Nevertheless, some detections of mid-IR H₂ emission from disks have been reported in the Herbig Ae/Be stars AB Aur (Bitner *et al.* 2007) and HD 97048 (Martin-Zaïdi *et al.* 2007), and in the T Tauri stars Sz 102, EC 74, EC 82, Ced 110 IRS6, EC 92, and ISO-Cha 237 (Lahuis *et al.* 2007). An interesting question to address is the reason for the high H₂ fluxes observed in these sources. A possibility is that an additional mechanism (X-rays, UV heating, e.g., Glassgold *et al.* 2007; Nomura *et al.* 2007) heats the molecular gas in the surface layer making $T_{\text{gas surf}} > T_{\text{dust surf}}$. An additional scenario is to invoke a gas-to-dust ratio much larger than the canonical value of 100 in the surface layer of the disk (e.g., as a result of dust coagulation and sedimentation). In order to explore the influence of a change in the gas-to-dust ratio and the thermal decoupling of gas

and dust in the surface layer, in Carmona *et al.* (2008) we calculated the expected 0–0 S(0) to 0–0 S(4) H₂ line fluxes from a M_{DISK} 0.01 M_⊙ disk model, as a function of the surface gas-to-dust ratio (ranging from 100 to 14000) for $T_{\text{gas surf}}/T_{\text{dust surf}}$ ranging from 1.0 to 2.0. We obtained that detectable 0–0 S(1) and 0–0 S(2) H₂ line flux levels can be achieved if $T_{\text{gas surf}}/T_{\text{dust surf}} > 2$ and if the gas-to-dust ratio in the surface layer is greater than 1000. *H₂ emission levels are very sensitive to departures from the thermal coupling between the molecular gas and dust in the surface layer.* Our results suggest that in the observed sources the molecular gas and the dust in the surface layer have not significantly departed from thermal coupling and that the gas-to-dust ratio in the surface layer is very likely lower than 1000. A definitive interpretation of our results awaits the development of future, more sophisticated models.

3. Searching for near-IR H₂ emission from protoplanetary disks with CRIRES

3.1. Motivation.

To probe the gas in the terrestrial planet forming region of the disk. H₂ near-IR lines probe hot gas at $T \sim 1000\text{--}3000$ K. This gas is located from a tenth of AU up to a few AU.

3.2. Previous work.

The $v=1\text{--}0$ S(1) H₂ line at 2.2218 μm has been detected in few classical T Tauri stars (CTTS): TW Hya, GG Tau A, LkCa 15 (Weintraub *et al.* 2000, Bary *et al.* 2002, 2003), AA Tau, CW Tau, UY Aur, GM Tau (Shukla *et al.* 2003), CS Cha (Weintraub *et al.* 2005), ECHAJ0843.3-7905 (Howat & Greaves 2007) and LkH α 264 (Itoh *et al.* 2003, Carmona *et al.* 2007), and in four weak-line T Tauri stars (WTTS): DoAr 21 (Bary *et al.* 2003), V773 Tau (Shukla *et al.* 2003), Sz33 and Sz 41 (Weintraub *et al.* 2005).

3.3. Observations & Data Reduction.

We observed the classical T Tauri star LkH α 264 and the debris disk 49 Cet, with the ESO-VLT cryogenic high-resolution ($R \sim 45000$, 6.6 km s^{-1}) infrared echelle spectrograph CRIRES (Käufl *et al.* 2004), mounted on ESO UT1 “Antu” 8-m telescope atop Cerro Paranal Chile, during the CRIRES science-verification phase (November 8 - 9, 2006). We employed the wave-ID 27/1/n and the wave-ID 25/-1/n, providing a spectral coverage from 2.0871 to 2.1339 μm and from 2.2002 to 2.2552 μm respectively. The observations were performed using a 46'' long, 0.4'' wide, north-south oriented slit, nodding the telescope 10'' along the slit. A random jitter smaller than 2'' was added to the telescope in addition to the nodding offset at each nodding position to correct for bad pixels and decrease systematics due to the detector. The total integration time was of 720s for LkH α 264 and of 240s for 49 Cet. Spectrophotometric standard stars at similar airmass to the science target were observed immediately following the science observations for performing the telluric correction. We searched for $v=1-0$ S(1) H₂ emission at 2.1218 μm , $v=1-0$ S(0) H₂ emission at 2.2233 μm and $v=2-1$ S(1) H₂ emission at 2.2477 μm . The data was reduced using the CRIRES pipeline and the ESO/CPL recipes. To correct for telluric absorption and flux-calibrate the science spectrum, the one-dimensional science spectrum obtained was divided by the one-dimensional spectrum of the standard star (for further details on the data reduction procedure, telluric correction and flux calibration see Sect. 2.1 of Carmona *et al.* 2007)

3.4. Results.

Our observations confirm the previous detections of the H₂ 1–0 S(1) line reported by Itoh *et al.* (2003). However, in contrast to Itoh *et al.* (2003), the H₂ 1–0 S(0) line is detected in our CRIRES spectra of LkH α 264. Our CRIRES observation show, for the first time, the simultaneous detection of the 1–0 S(1) and 1–0 S(0) H₂ line from a protoplanetary disk.

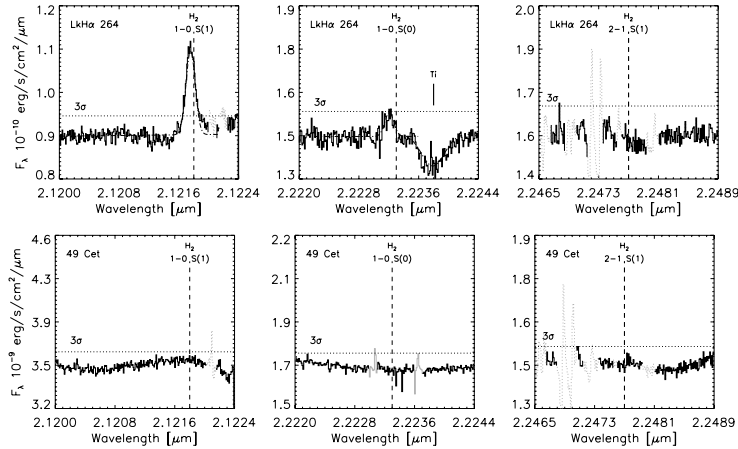


Figure 3. CRIRES spectra of LkH α 264 (upper panels) and 49 Cet (lower panels) in the regions of the H $_2$ $v=1-0$ S(1), H $_2$ $v=1-0$ S(0) and H $_2$ $v=2-1$ S(1) emission lines. The H $_2$ $v=1-0$ S(1) and the H $_2$ $v=1-0$ S(0) lines are detected in LkH α 264. Photospheric Ti features at 2.2217 (not shown) and 2.2238 μm (central upper panel) are observed in LkH α 264. The Gaussian fits to the detected lines are illustrated in dash-dot lines. The H $_2$ $v=2-1$ S(1) line is not present in LkH α 264. In the case of 49 Cet none of the three H $_2$ features are present in emission or absorption. Horizontal dotted lines show the 3σ continuum flux limits. The spectra are not corrected for V_{LSR} of the star. Regions of poor telluric correction are in gray-dotted lines in the spectra (Carmona *et al.* 2007).

Table 2. H $_2$ near-IR emission line fluxes and upper limits

Star	H $_2$ line	λ [μm]	continuum [$\times 10^{-10}$ ergs s $^{-1}$ cm $^{-2}$ μm^{-1}]	line flux ^a [$\times 10^{-14}$ ergs s $^{-1}$ cm $^{-2}$]
LkH α 264	1-0 S(1)	2.1218	0.9 (0.04)	0.30
	1-0 S(0)	2.2233	1.5 (0.08)	0.10
	2-1 S(1)	2.2477	1.6 (0.10)	<0.05
49 Cet	1-0 S(1)	2.1218	35 (1.0)	<0.5
	1-0 S(0)	2.2233	17 (1.8)	<0.9
	2-1 S(1)	2.2477	15 (3.1)	<1.6

Notes: ^a Upper limits calculated using a line $FWHM$ of 6.6 km s $^{-1}$.

The H $_2$ 2-1 S(1) line is not seen in LkH α 264. The central wavelength of the 1-0 S(1) H $_2$ emission in LkH α 264 was measured to be 2.121757 ± 0.000005 μm . This corresponds to a velocity shift of -5.6 ± 1.0 km s $^{-1}$ (taking into account Earth's velocity at the time of observation), a velocity coincident with the rest velocity of the star (-5.9 ± 1.2 km s $^{-1}$, Itoh *et al.* 2003). The H $_2$ 1-0 S(0) feature at 2.2233 μm is detected with a 3σ level confidence. Employing a Gaussian fit, the central wavelength of the line found is 2.22321 ± 0.00005 μm . This corresponds to a velocity shift of -12 ± 7 km s $^{-1}$ which is in agreement with the velocity shift found in the 1-0 S(1) line. The $FWHM$ of the lines are 20.6 ± 1 km s $^{-1}$ for the H $_2$ 1-0 S(1) line and 19.8 ± 1 km s $^{-1}$ for the H $_2$ 1-0 S(0) line. The measured line fluxes are 3.0×10^{-15} and 1.0×10^{-15} erg s $^{-1}$ cm $^{-2}$ for the H $_2$ 1-0 S(1) and H $_2$ 1-0 S(0) respectively. The 2-1 S(1) H $_2$ line is not observed in LkH α 264. Assuming a $FWHM$ of 6.6 km s $^{-1}$ (CRIRES resolution), a 3σ flux upper limit of 5.3×10^{-16} ergs s $^{-1}$ cm $^{-2}$ is derived for the line. Assuming an error of 20% in the flux calibration of the spectra, the 1-0 S(0)/1-0 S(1) line ratio in LkH α 264 is 0.33 ± 0.1 and the 2-1 S(1)/1-0 S(1) line ratio is <0.2 . These line ratios are consistent with the line ratios of a gas at LTE at a temperature cooler than 1500 K (Mouri 1994).

In the case of 49 Cet none of the three H $_2$ features are present in emission or absorption. A summary of the 3σ flux upper limits is presented in Table 2.

3.5. Discussion.

H₂ emission in LkH α 264: a disk or an outflow? The first question is to determine whether the H₂ emission observed in LkH α 264 originates in an outflow (shock excited emission) or in a disk. The small velocity shift, the line shape (well reproduced by a disk model), and the fact that the emission is spatially unresolved are not in favor of shock excited H₂. An additional strong argument against shock excitation of H₂ is that LkH α 264 does not exhibit [OI] forbidden emission at 6300 Å (Cohen and Kuhi 1979); a classical signature of outflows in T Tauri stars. The lack of this line indicates that in LkH α 264 the outflow is not present or at least that it is very weak. We conclude that the H₂ emission observed in LkH α 264 originates very likely in a disk.

Excitation mechanism of the H₂ emission in LkH α 264. Thermal and non-thermal excitation mechanisms are distinguishable on the basis of line ratios (Mouri 1994 and references therein). With Figure 3b of Mouri (1994), we find that the measured 1–0 S(0)/1–0 S(1) (0.33 ± 0.1) and the 2–1 S(1)/1–0 S(1) (< 0.2) line ratios in LkH α 264 are consistent with thermal emission of a gas cooler than 1500 K thermally excited by UV photons (i.e., UV photons NOT in the Lyman-Werner band at 912–1108 Å). Assuming a gaussian error distribution, given the line 1–0 S(0)/1–0 S(1) ratio measured, the probability that the heating mechanism is X-ray excitation is less than 1%.

H₂ emitting region and inclination of the LkH α 264 disk. The spectral resolution of CRILES ($\approx 6.6 \text{ km s}^{-1}$) and the thermal width of a 1500 K line ($\approx 2.4 \text{ km s}^{-1}$) are significantly smaller than the *FWHM* of 20 km s^{-1} of the H₂ lines observed in LkH α 264. Therefore, the line width must be linked to the dynamics of the gas in the region that is emitting the line. Implementing the two-layer Chiang and Goldreich (1997) disk model code CG plus (Dullemond *et al.* 2001), we modeled the disk around LkH α 264, and found that the regions of the disk with a surface layer at $T_s < 1500 \text{ K}$ are located at $R > 0.1 \text{ AU}$. Modeling of the single peaked 1–0 S(1) line shape (for the details of the model see Carmona *et al.* 2007) indicates that the disk is close to face-on ($i < 35^\circ$). The best model fit suggests that the disk of LkH α 264 is inclined 20° for a H₂ emitting region extending from 0.1 to 10 AU with a power law relation of the intensity as a function of radius with exponent $\alpha = -2$. If the 1–0 H₂ S(1) line intensity decreases with an exponent $\alpha = -2$ as a function of radius, then 50% of the line flux is produced within 0.1 AU and 1 AU of the LkH α 264 disk, 40% of the line flux is emitted within 1 and 7 AU and the rest of the flux at larger radii.

Mass of optically thin hot H₂ in LkH α 264 and 49 Cet. Using a similar equation to Eq. 2.1 from the 1–0 S(1) H₂ line flux in LkH α 264 and the upper limit to the flux of the same line in 49 Cet, assuming optically thin gas in LTE, we calculated the mass H₂ gas at 1500 K in both sources. We found that there are a few lunar masses of hot H₂ gas in the inner disk of LkH α 264 and less than a tenth of a lunar mass of hot H₂ in the inner disk of 49 Cet. The lack of H₂ ro-vibrational emission in the spectra of 49 Cet, combined with non detection of pure rotational lines of H₂ (Chen *et al.* 2006) and the absence of H α emission suggest that the gas in the inner disk of 49 Cet has dissipated. These results together with the previous detection of ¹²CO emission at sub-mm wavelengths (Zuckerman *et al.* 1995; Dent *et al.* 2005) point out that the disk of 49 Cet should have a large inner hole, and it is strongly suggestive of theoretical scenarios in which the disk disappears inside-out. This could be due to inside-out photoevaporation, or to the presence of an unseen low-mass companion(s).

Detections and non detections of near-IR H₂ emission from T Tauri stars. In Carmona *et al.* (2007) we present a detailed comparative analysis of the physical properties of classical T Tauri stars in which the H₂ 1–0 S(1) line has been detected versus

non-detected. We found that there is a higher chance of observing the H₂ near-infrared lines in CTTS with a high $U - V$ excess and a strong H α line. This result suggests that there is a higher probability of detecting the H₂ 1–0 S(1) line in systems with high accretion. In contrast to weak-lined T Tauri stars, there is no apparent correlation between the X-ray luminosity and the detectability of the H₂ 1–0 S(1) line in classical T Tauri stars. Taken as a group, LkH α 264 and the CTTS in which near-IR H₂ emission has been detected exhibit typical properties of classical T Tauri stars. Therefore, we expect near-IR ro-vibrational H₂ lines from T Tauri disks to be detected on a routine basis in the near future.

4. Conclusions

The advent of infrared high-resolution spectrographs mounted on large aperture telescopes allow for the first time the study of the gaseous component of the disk in the planet forming region ($R < 50$ AU). Here, we presented two efforts for constraining the properties of the molecular hydrogen in a sample of Herbig Ae/Be, T Tauri stars and a debris disk.

In the first project, we searched for fundamental emission of H₂ at 17 and 12 micron employing VISIR, ESO's mid-infrared high resolution spectrograph. Our aim was to probe the warm gas ($T \sim 150 - 1000$ K) in the giant planet forming region of the disk. None of the sources show fundamental H₂ emission. From the upper limits on the line fluxes derived, we estimated that there is less than a tenth of Jupiter mass of optically thin H₂ at $T = 150$ K and less than a few Earth mass at higher temperatures. These results are consistent with the two-layer disk model (CG 97). In this model only the molecular gas in the optically thin surface layer of the disk emits the H₂ lines. Typical Herbig Ae disks have a very small mass of molecular gas ($< 10^{-2} M_J$) in the surface layer. We calculated the expected H₂ fluxes assuming a gas in LTE and a gas-to-dust ratio of 100. We found that the line fluxes are of the order of 10^{-16} to 10^{-17} erg s⁻¹ cm⁻², two orders of magnitude fainter than our detection limits. If the two-layer model is correct we are “blind” to the H₂ gas in the interior layer. This result explains the numerous non-detections of mid-IR H₂ emission towards several pre-main sequence stars with disks. The few detections reported can be explained as the result of a thermal decoupling of the gas and the dust (gas hotter than the dust) or a change in gas-to-dust ratio in the surface layer. Our non-detections suggest that in our sources the H₂ and dust in the surface layer have not significantly departed from thermal coupling ($T_{\text{gas}}/T_{\text{dust}} < 2$) and the gas-to-dust ratio in the surface layer is very likely < 1000 .

In the second project, we searched for near-infrared ro-vibrational H₂ emission from LkH α 264 a classical T Tauri star and 49 Cet a debris disk with detections of CO emission in the sub-mm. Our aim was to probe the hot gas ($T \sim 1000$ K) present in the terrestrial planet region of the disks. We detected the 1–0 S(1) and 1–0 S(0) H₂ lines in LkH α 264 and derived an upper limit to the 2–1 S(1) H₂ line. Given that the velocity of the lines is coincident with the velocity of LkH α 264, that both lines have a similar *FWHM* of 20 km s⁻¹ (well reproduced by a disk model) and that LkH α 264 does not have evidence for outflows, we concluded that the observed emission very likely arises from a circumstellar disk. This is the first time that the 1–0 S(1) and 1–0 S(0) H₂ lines are detected in a protoplanetary disk. From the line ratios, we deduced that the H₂ gas emitting the line is at temperature lower than 1500 K and that the gas is likely thermally excited by UV photons. Modeling of the disk around LkH α 264 and the shape of the 1–0 S(1) H₂ line suggests that the disk is observed close to face on ($i < 35^\circ$) and that most of the near-IR H₂ emission is produced at less than 1 AU. The flux of the 1–0 S(1) H₂ line indicates that there is a few lunar masses of hot ($T \sim 1500$ K) optically thin H₂ in the inner disk

($R \sim 0.1$ AU) of LkH α 264. In the case of 49 Cet none of the three H₂ near-IR lines are observed in the spectra. The flux upper limits on the 1-0 S(1) H₂ line indicate that there is less than a fraction of lunar mass in the inner ($R < 1$ AU) disk of 49 Cet. This result combined with the absence of near-IR excess in the spectral energy distribution, the lack of H α in emission and the non-detection of mid-IR H₂ emission (from Spitzer), suggest that the inner disk of 49 Cet has a hole in the gas and in the dust.

These projects show how high-resolution infrared spectroscopy is a useful tool for constraining the structure of protoplanetary disks, even in the case non-detections. Future studies, in larger samples are essential to identify statistical trends between the gas and dust properties of the disk as a function of the age and multiplicity.

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